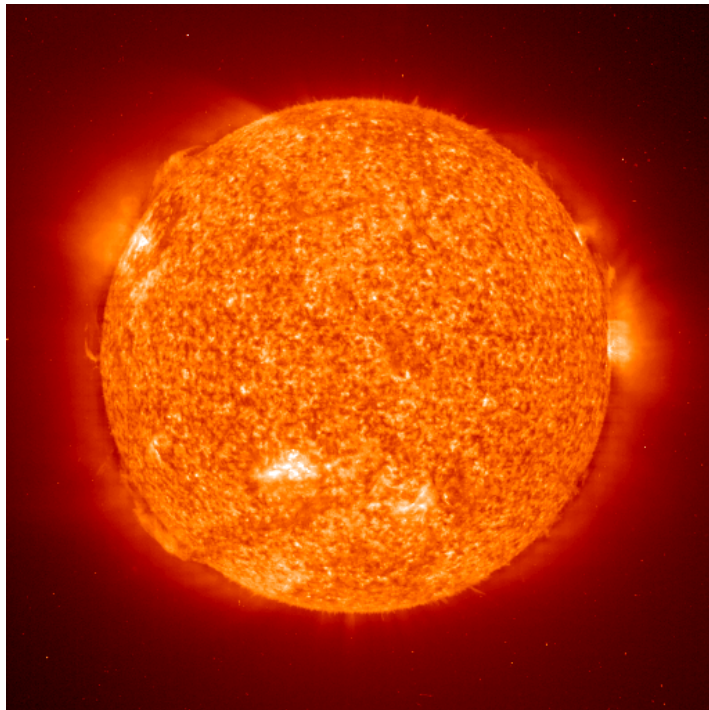


Tokamak Operational Scenarios

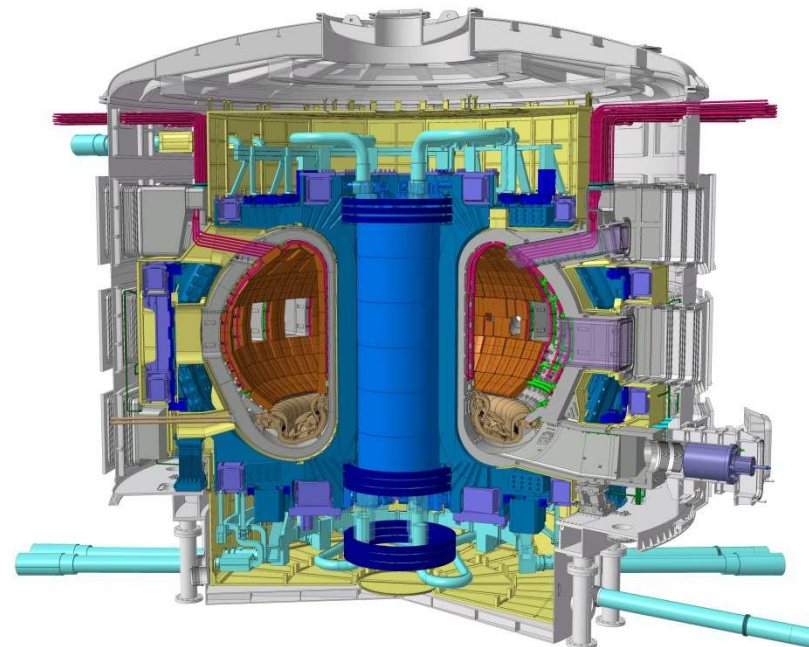


DPG Advanced Physics School
,The Physics of ITER'
Bad Honnef, 25.09.2014

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- What is a 'tokamak (operational) scenario'?
- Optimisation strategies for tokamak scenarios
- Conventional scenarios
- Advanced scenarios
- Summary and conclusions



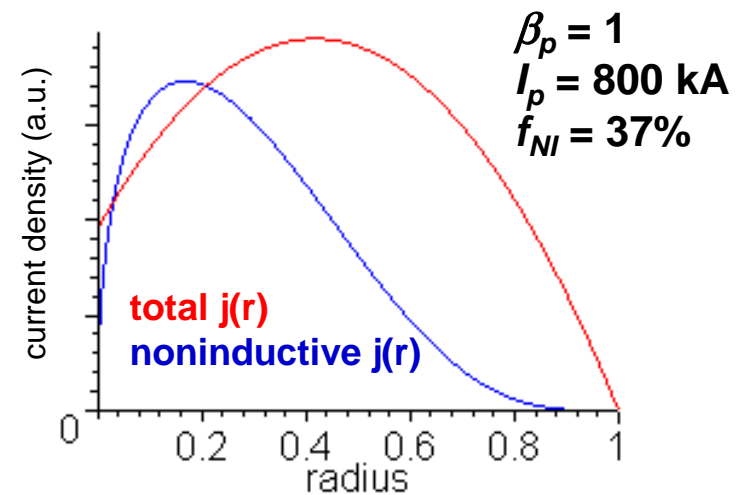
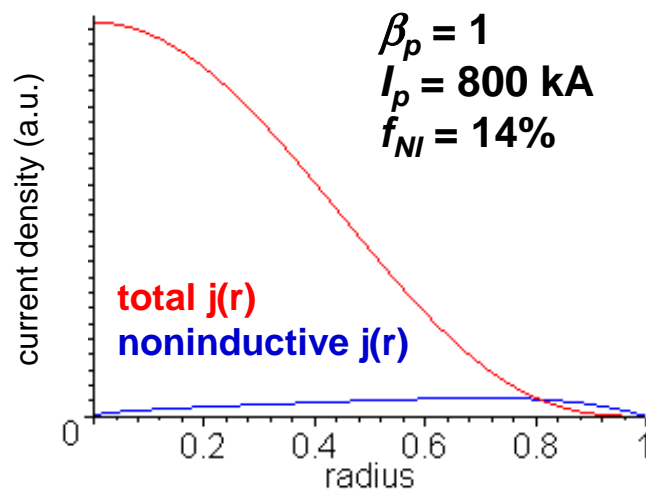
What is a 'tokamak scenario'?



A tokamak (operational) scenario is a recipe to run a tokamak discharge

Plasma discharge characterised by

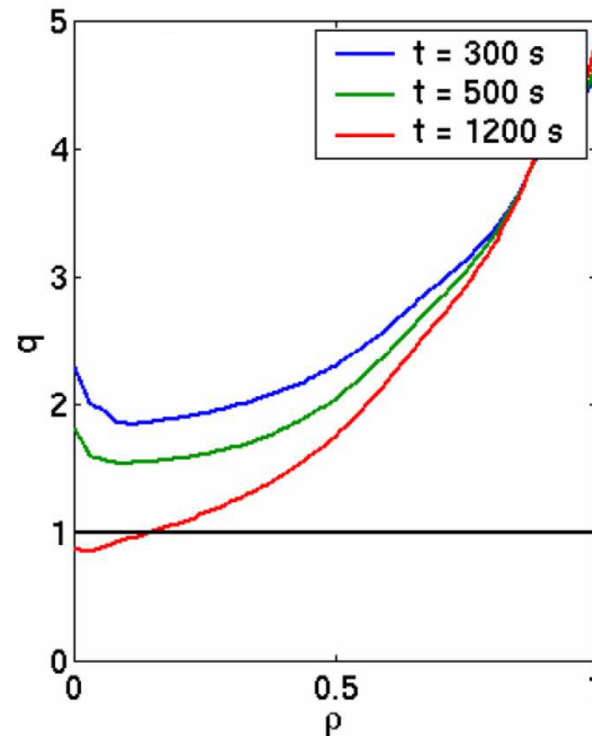
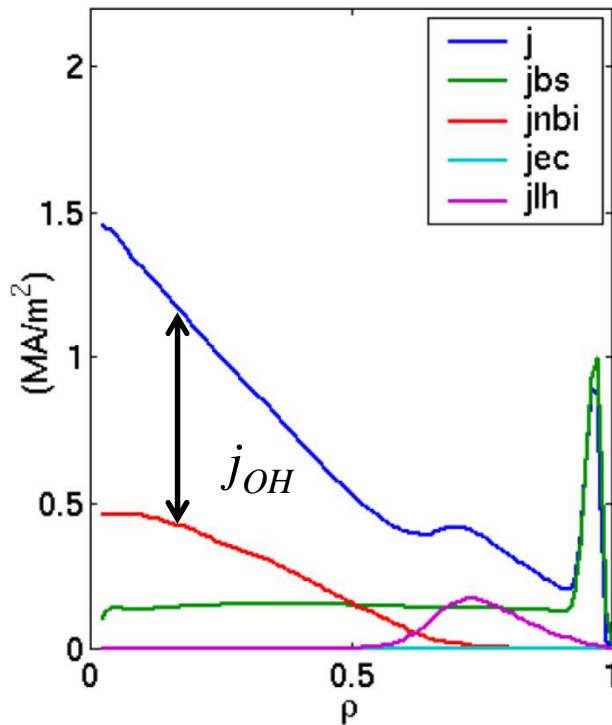
- external control parameters: B_t , R_0 , a , κ , δ , P_{heat} , Φ_D ...
- integral plasma parameters: $\beta = 2\mu_0 \langle p \rangle / B^2$, $I_p = 2\pi \int j(r) r dr$...
- plasma profiles: pressure $p(r) = n(r) \cdot T(r)$, current density $j(r)$



→ operational scenario best characterised by *shape* of $p(r)$, $j(r)$



Control of the profiles $j(r)$ and $p(r)$ is limited



safety factor:

$$q \approx \frac{r}{R} \frac{B_{tor}}{B_{pol}} \propto \frac{r^2}{R} \frac{B_{tor}}{I_p}$$

ITER Q=10 simulation

- ohmic current coupled to temperature profile via $\sigma \sim T^{3/2}$
→ inductive current profiles always peaked, q -profiles monotonic
- external heating systems drive current, but with limited efficiency (typically less than 0.1 A per 1 W under relevant conditions)...
- pressure gradient drives toroidal 'bootstrap' current: $j_{bs} \sim (r/R)^{1/2} \nabla p / B_{pol}$



Control of the profiles $j(r)$ and $p(r)$ is limited



Pressure profile determined by combination of heating / fuelling profile and radial transport coefficients

- ohmic heating coupled to temperature profile via $\sigma \sim T^{3/2}$
- external heating methods allow for some variation – ICRH/ECRH deposition determined by B -field, NBI has usually broad profile
- under reactor-like conditions, dominant α -heating $\sim (nT)^2$
- gas puff is peripheral source of particles, pellets further inside

Shape of profile will strongly depend on (turbulent) heat conductivity and particle diffusivity (i.e. be 'self-organised')



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Figure of merit for fusion performance $nT\tau$

IPP

Power P_{loss} needed to sustain plasma

- determined by thermal insulation:

$$\tau_E = W_{plasma} / P_{loss} \text{ ('energy confinement time')}$$

Fusion power increases with W_{plasma}

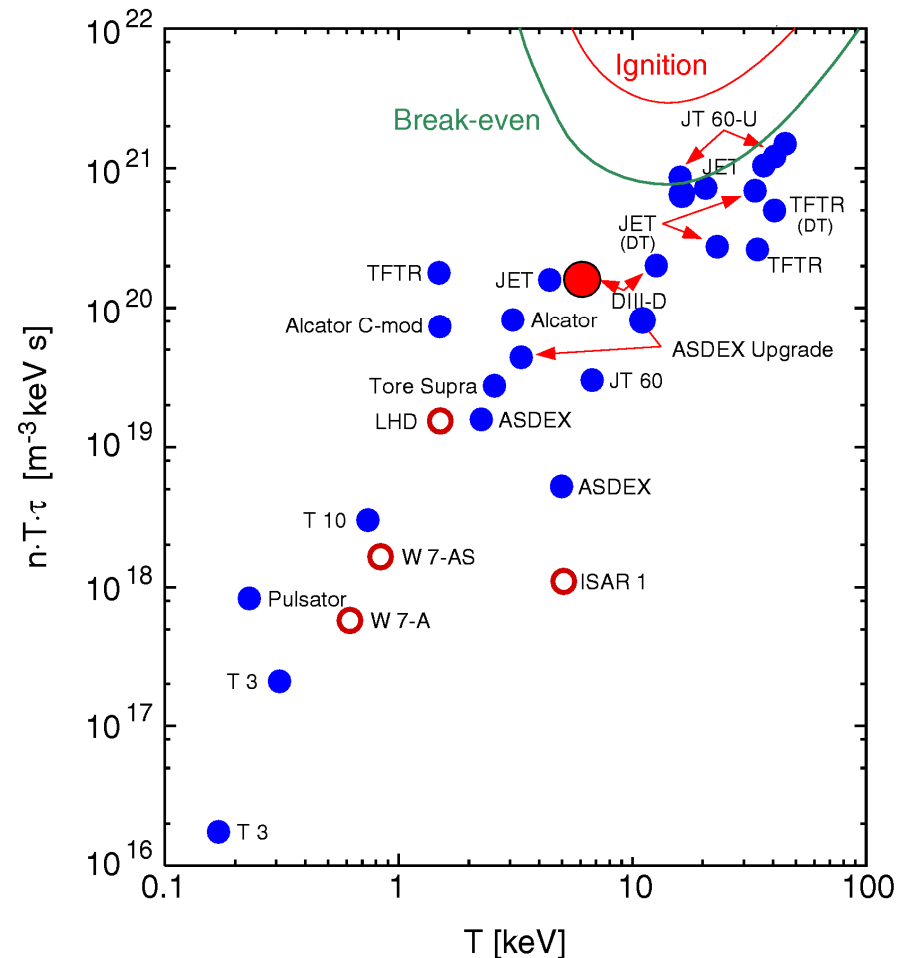
- $P_{fus} \sim n_D n_T \langle \sigma v \rangle \sim n_e^2 T^2 \sim W_{plasma}^2$

Present day experiments: P_{loss} compensated by external heating

- $Q = P_{fus} / P_{ext} \approx P_{fus} / P_{loss} \sim nT\tau_E$

Reactor: P_{loss} compensated by α -(self)heating

- $Q = P_{fus} / P_{ext} = P_{fus} / (P_{loss} - P_{\alpha}) \rightarrow \infty$ (ignited plasma)





Optimisation of $nT\tau_E$: ideal pressure limit

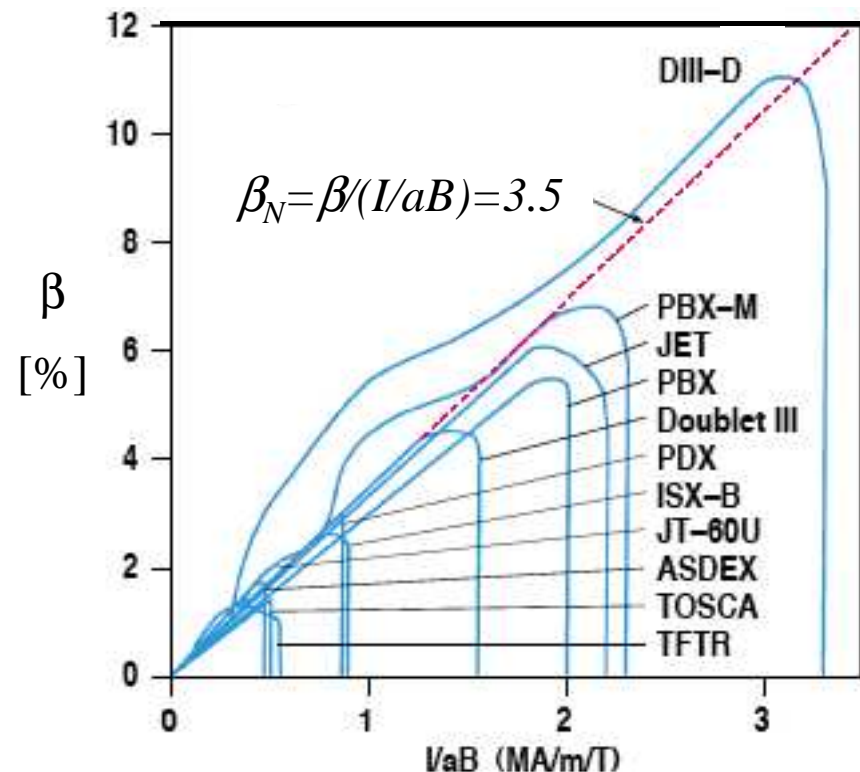


Optimising nT means high pressure and, for given magnetic field, high $\beta = 2\mu_0 \langle p \rangle / B^2$

This quantity is limited by magneto-hydrodynamic (MHD) instabilities

‘Ideal’ MHD limit (ultimate limit, plasma unstable on Alfvén time scale $\sim 10 \mu\text{s}$, only limited by inertia)

- ‘Troyon’ limit $\beta_{max} \sim I_p/(aB)$, leads to definition of $\beta_N = \beta/(I_p/(aB))$





Optimisation of $nT\tau_E$: resistive pressure limit

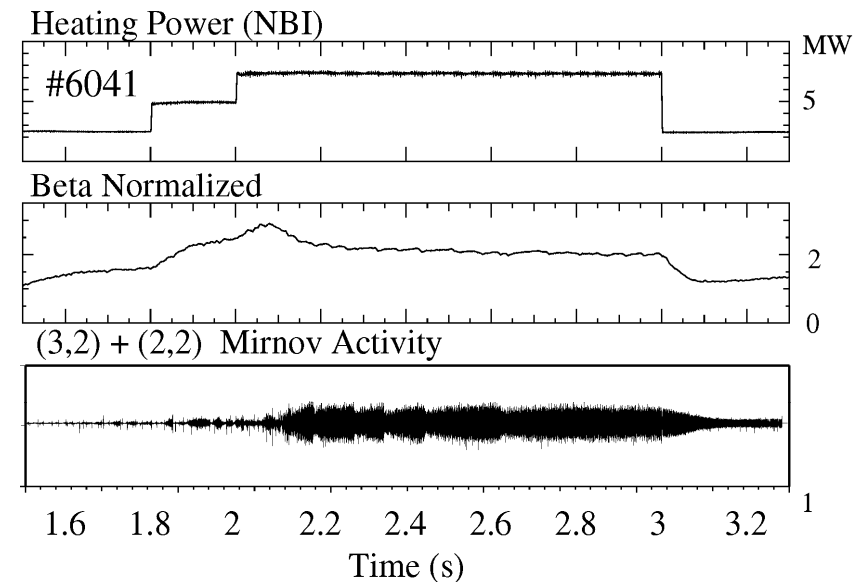


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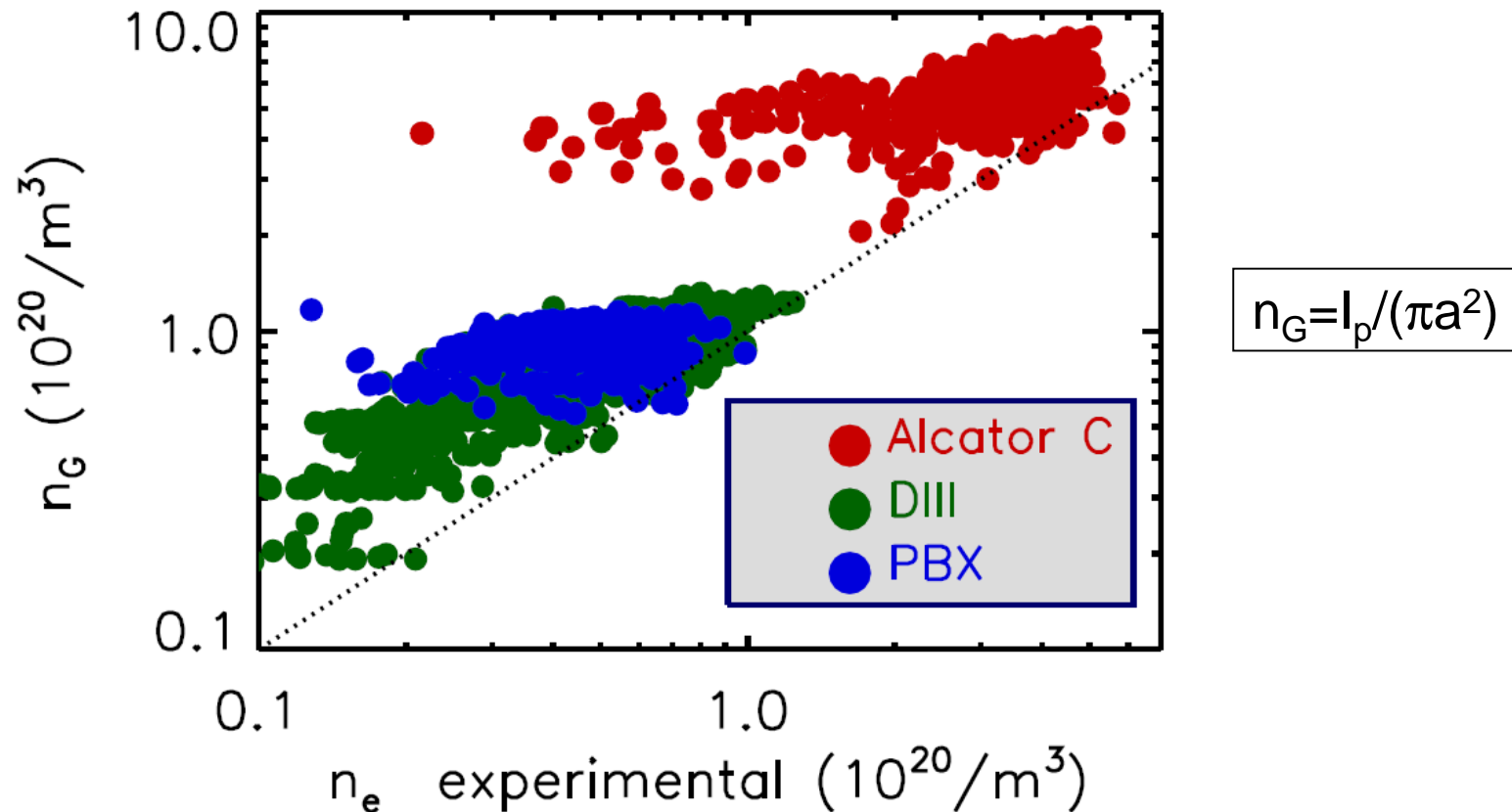
‘Resistive’ MHD limit (on local current redistribution time scale ~ 100 ms)

Neoclassical Tearing Mode (NTM)





Optimisation of $nT\tau_E$: density limit



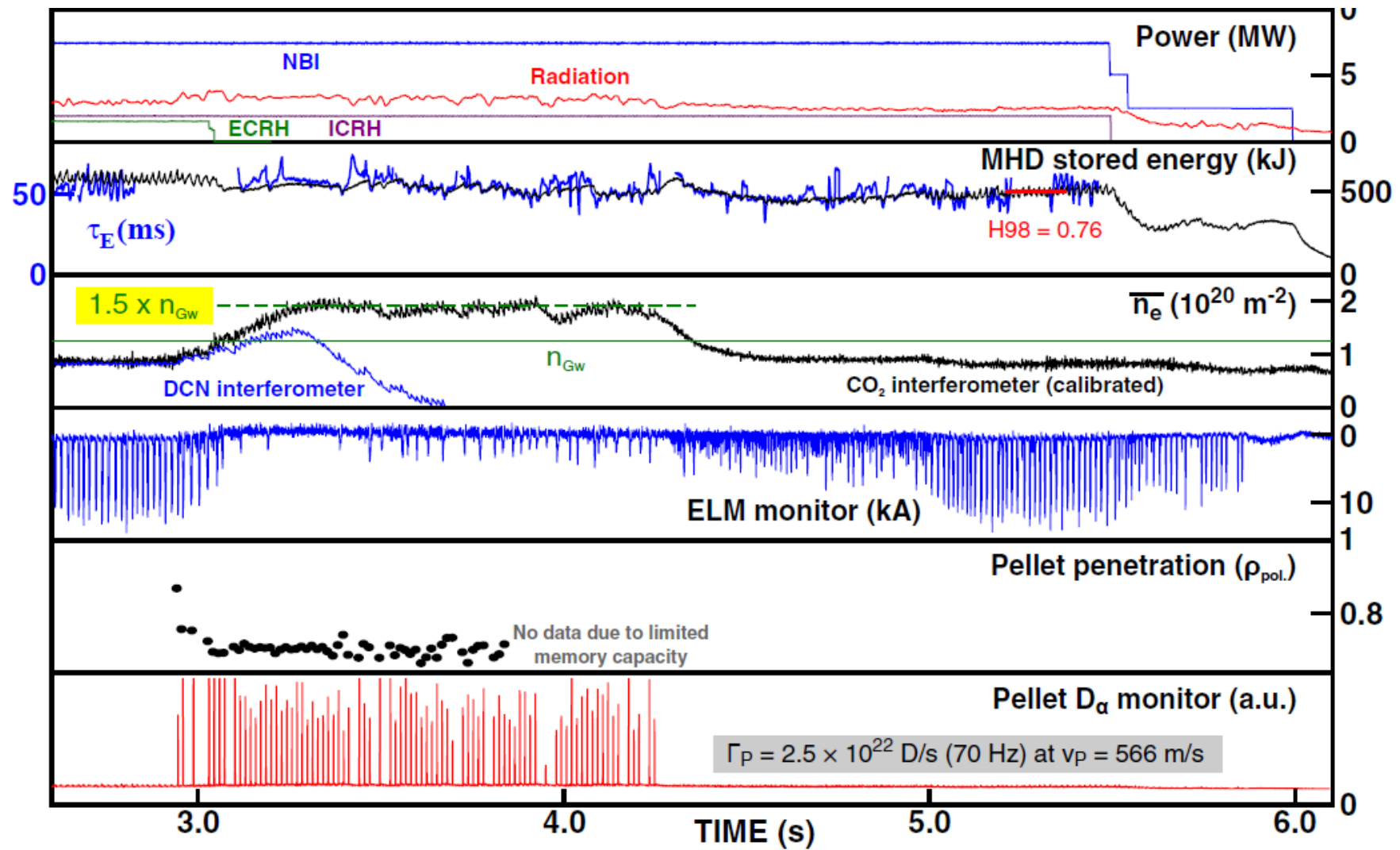
Empirical 'Greenwald-limit' describes well maximum density

- seems to be linked to a change in edge transport at $n \sim n_G$
- can be overcome if density profile shape is varied (peaked)



Operation at $n/n_{GW} > 1$

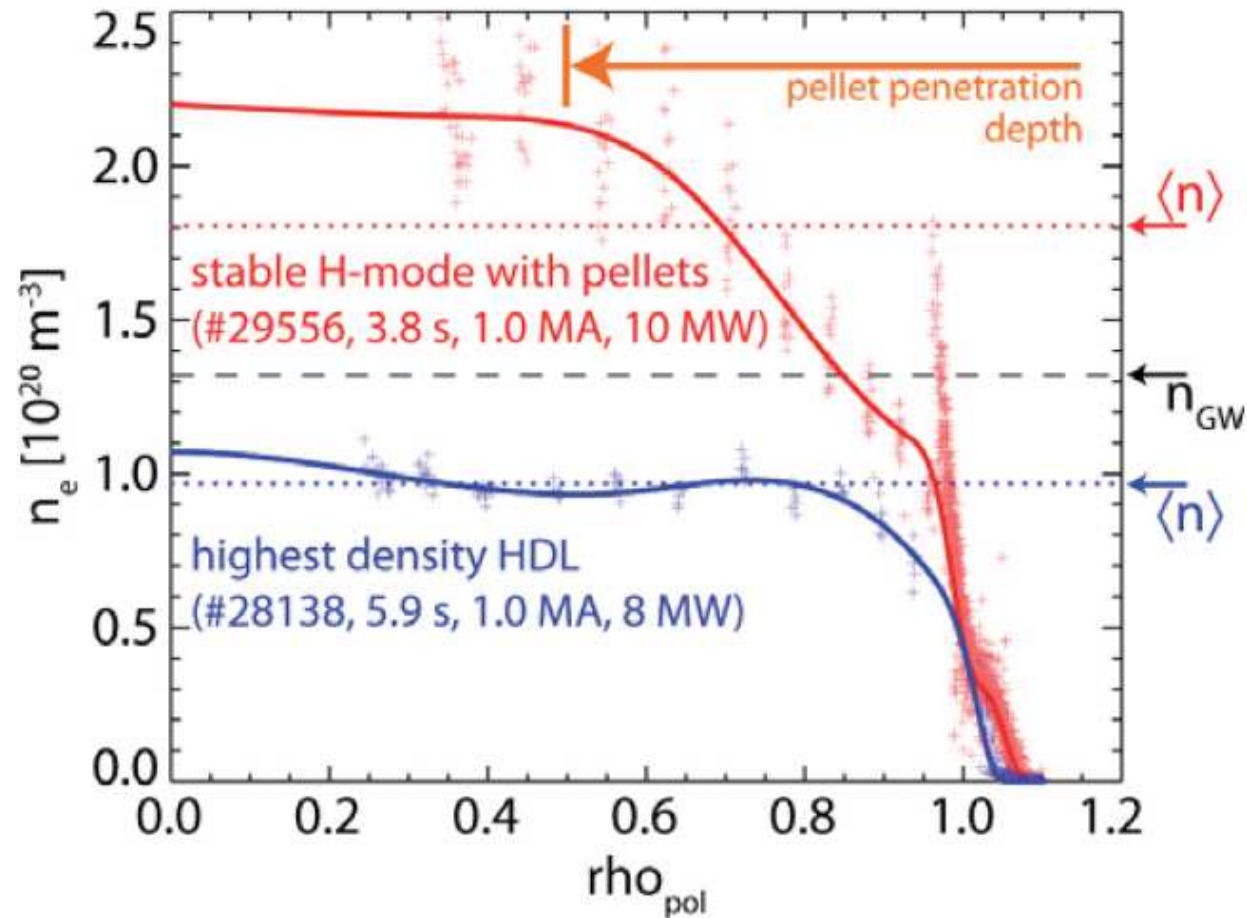
IPP



Stable operation at $n/n_{GW} = 1.5$ using pellets



Operation at $n/n_{\text{GW}} > 1$

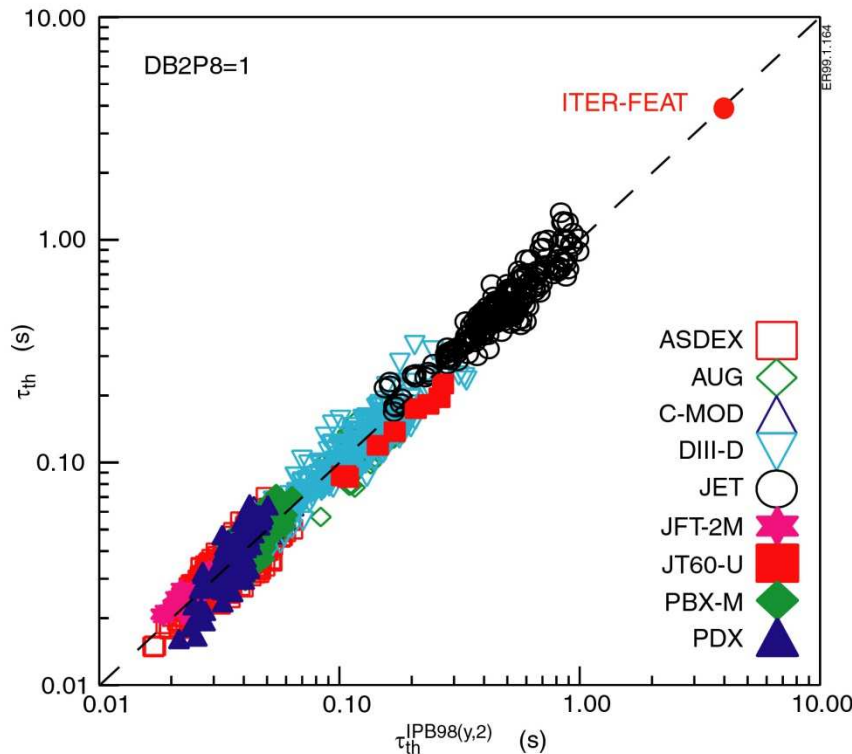


Edge density stays below n_{GW} in all cases (up to $n_{e0} = 4 \times n_{\text{GW}}!$)

- H-mode density limit = edge density limit



Optimisation of $nT\tau_E$: confinement scaling



Empirical ITER 98(p,y) scaling:

$$\tau_E \sim H I_p^{0.93} P_{heat}^{-0.63} B_t^{0.15} \dots$$

Empirical confinement scalings show linear increase of τ_E with I_p

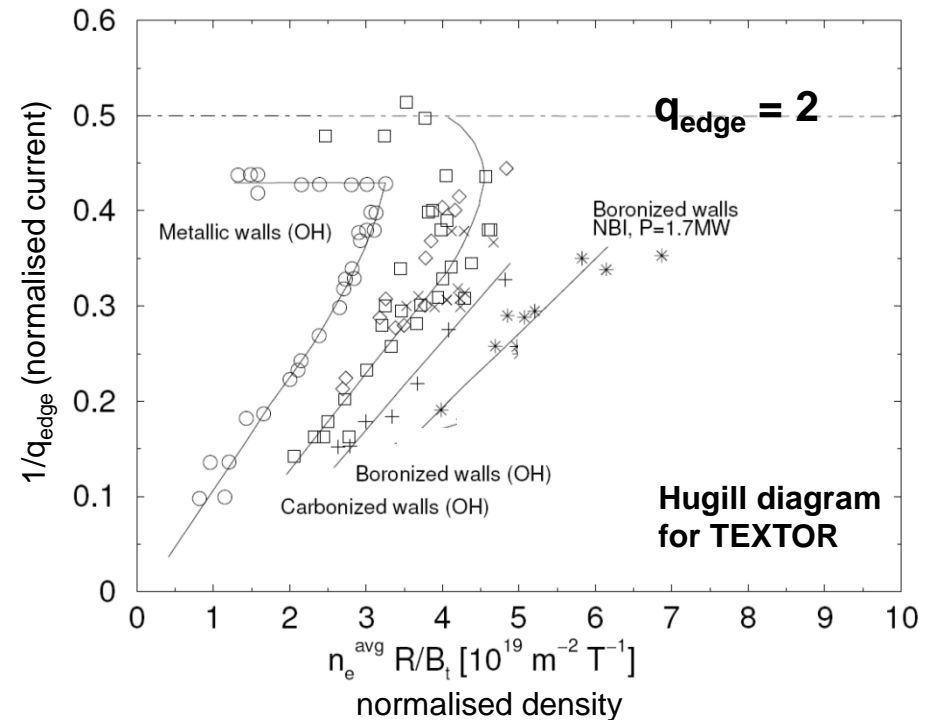
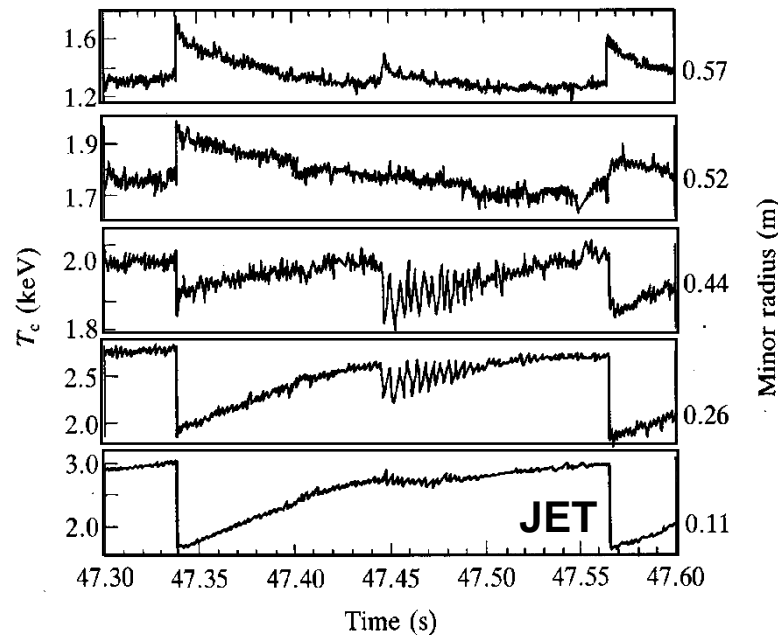
- note the power degradation (τ_E decreases with P_{heat} !)
- 'H-factor' H measures the quality of confinement relative to the scaling



Optimisation of $nT\tau_E$: current limit



BUT: for given B_t , I_p is limited by current gradient driven MHD instabilities

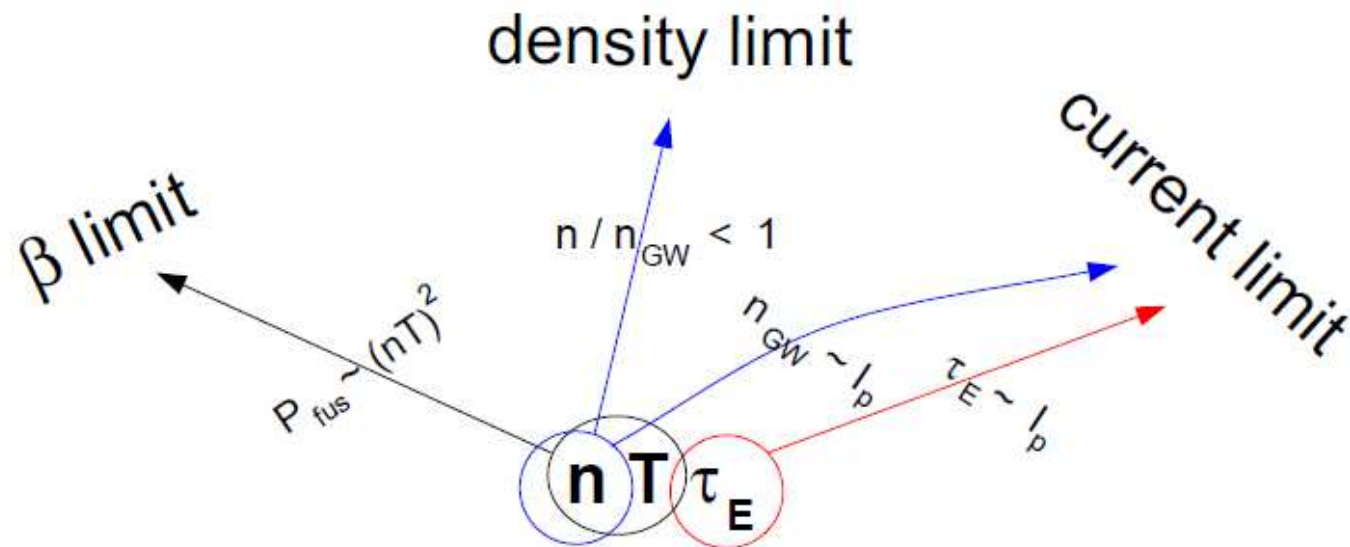


Limit to safety factor $q \sim (r/R) (B_{tor}/B_{pol})$

- for $q < 1$, tokamak unconditionally unstable \rightarrow central 'sawtooth' instability
- for $q_{edge} \rightarrow 2$, plasma tends to disrupt (external kink) – limits value of I_p



Optimisation of $nT\tau_E$



Optimising for $Q = P_{\text{fus}} / P_{\text{ext}}$ drives operational point close to operational limits



For steady state tokamak operation, high I_p is not desirable:

Tokamak operation without transformer: current 100% noninductive

- external CD has low efficiency (remember less than 0.1 A per W)
- internal bootstrap current high for high $j_{bs} \sim (r/R)^{1/2} \nabla p / B_{pol}$

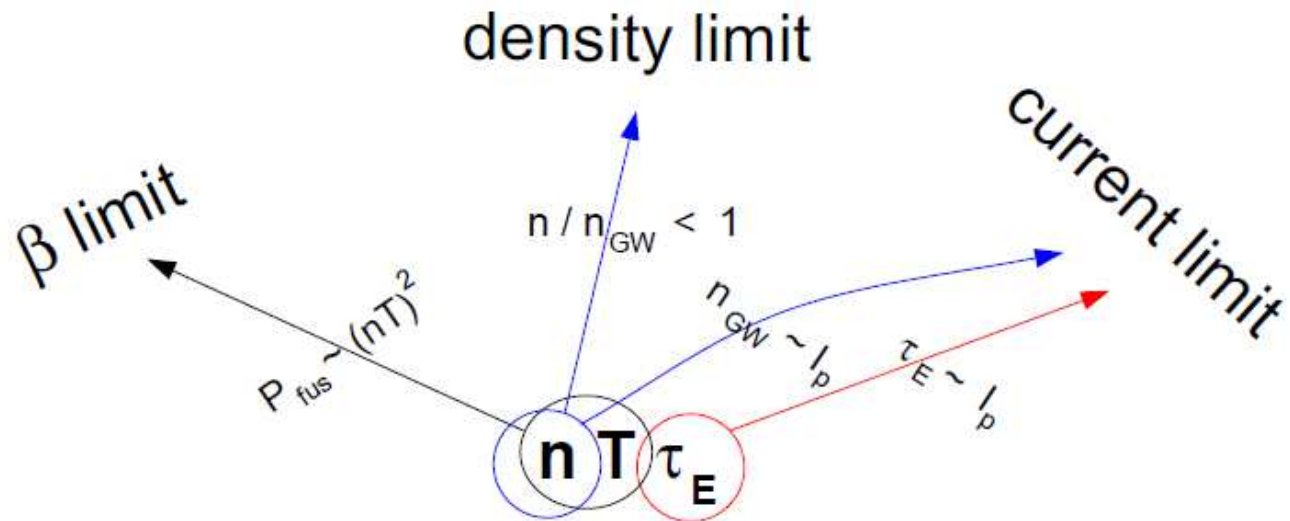
$$\rightarrow f_{NI} \sim I_{bs} / I_p \sim p / B_{pol}^2 \sim \beta_{pol}$$

‘Advanced’ scenarios, which aim at steady state, need high β , low I_p , have to make up for loss in τ_E (e.g. through transport barriers)

Without the ‘steady state’ boundary condition, a tokamak scenario is called ‘conventional’



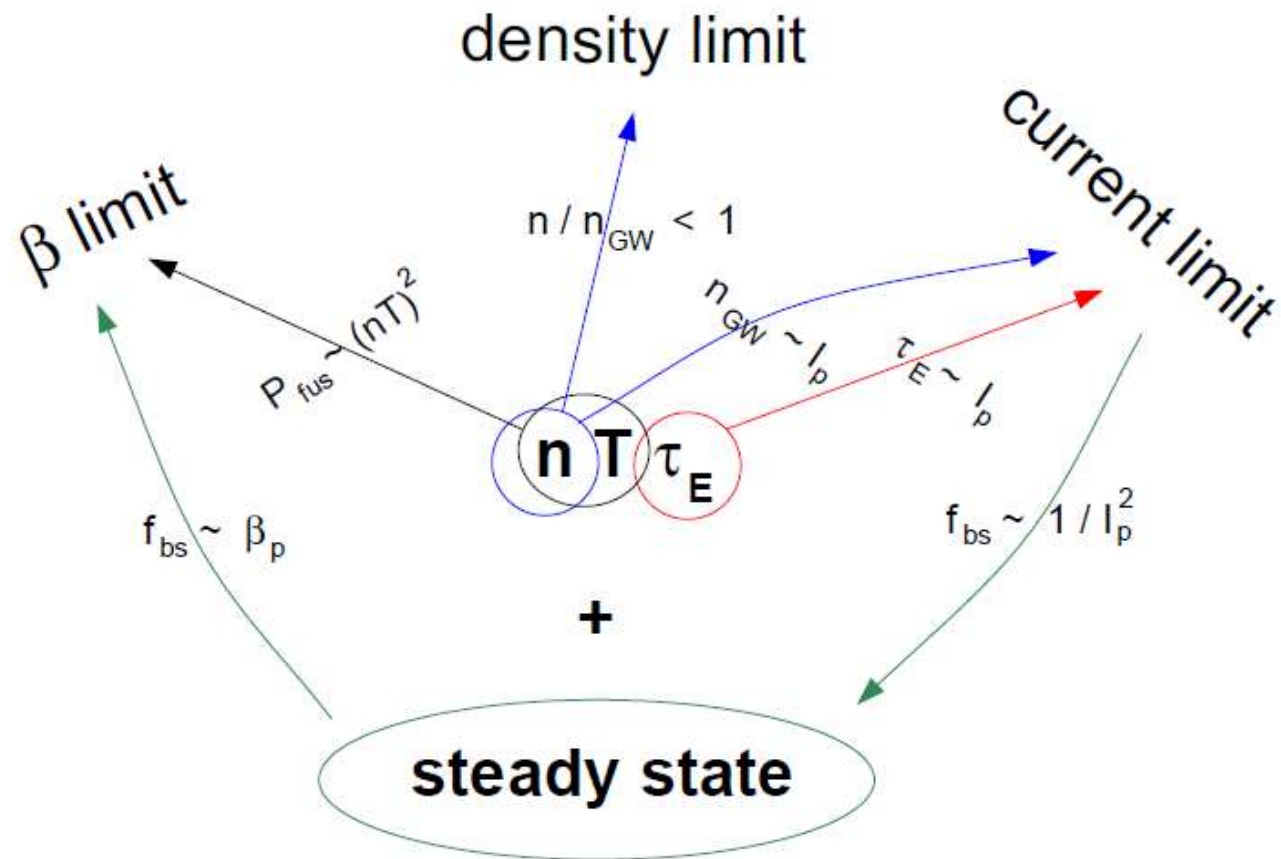
Optimisation of $nT\tau_E$



Optimising for $Q = P_{\text{fus}} / P_{\text{ext}}$ drives operational point close to operational limits



Optimisation of $nT\tau_E$



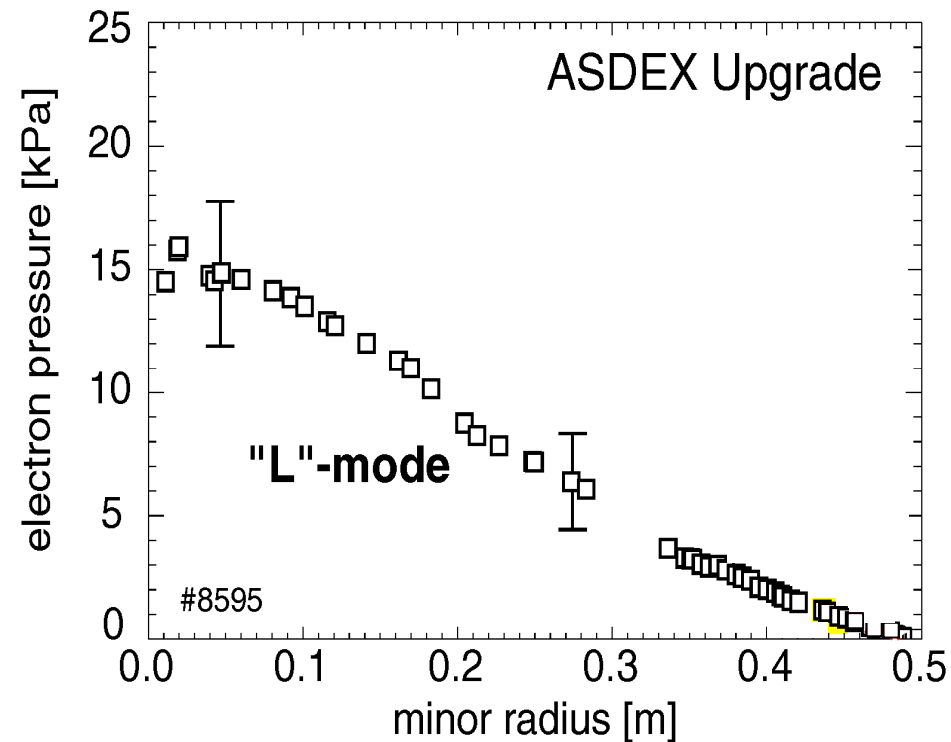
Including the steady state constraint emphasizes the β -limit
(and de-emphasizes current limit)



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The (low confinement) L-mode scenario

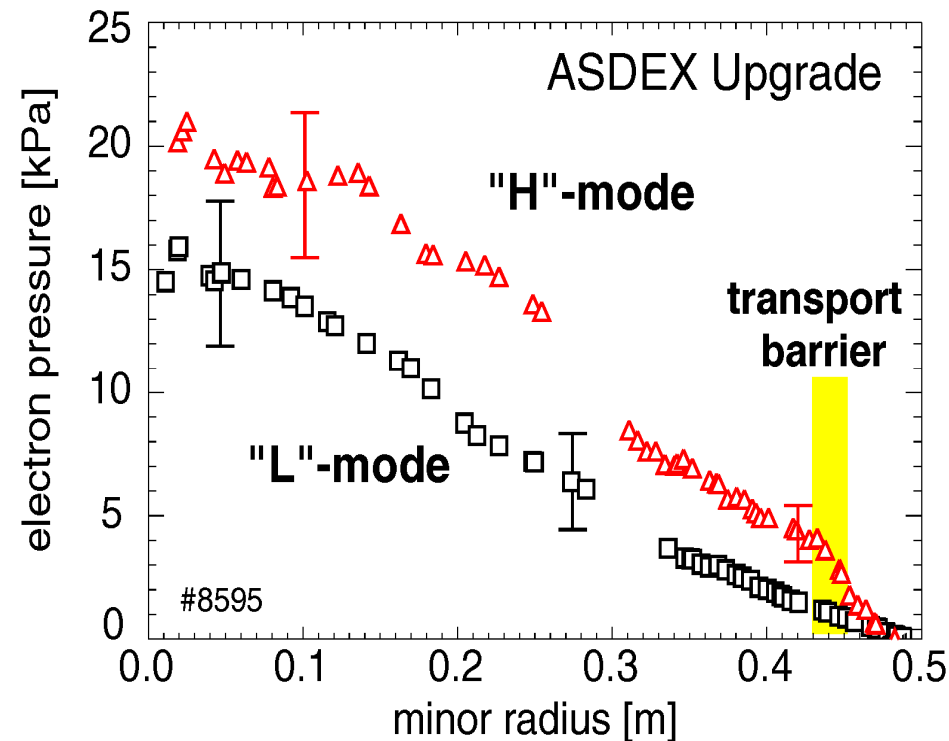


Standard scenario without special tailoring of geometry or profiles

- central current density usually limited by sawteeth
- temperature gradient sits at critical value over most of profile
- extrapolates to very large ($R > 10$ m, $I_p > 30$ MA) pulsed reactor



The (high confinement) H-mode scenario

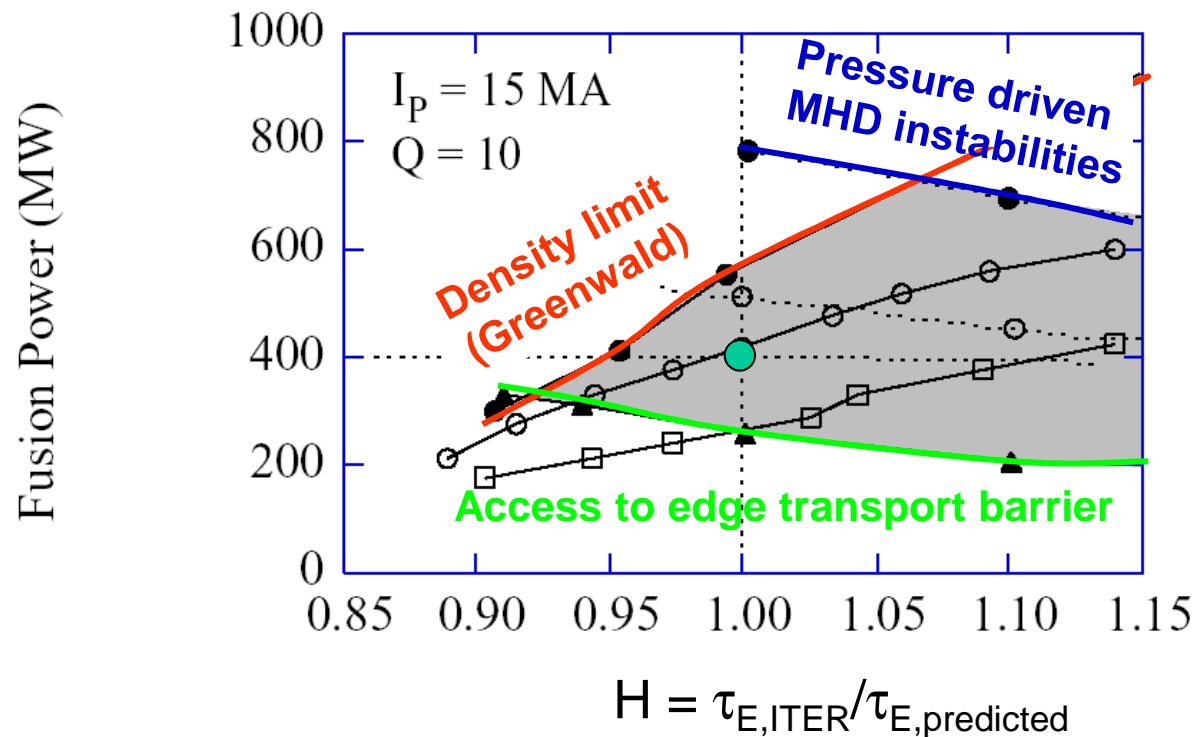


With hot (low collisionality) conditions, edge transport barrier develops

- gives higher boundary condition for 'stiff' temperature profiles
- global confinement τ_E roughly factor 2 better than L-mode
- extrapolates to more attractive ($R \sim 8$ m, $I_p \sim 20$ MA) pulsed reactor



H-mode is ITER standard scenario for $Q=10$...



The design point allows for...

- ...achieving $Q=10$ with conservative assumptions
- ...incorporation of 'moderate surprises'
- ...achieving ignition ($Q \rightarrow \infty$) if surprises are positive



...but some open issues remain...



Need to minimise ELM impact on divertor

- reduce power flow to divertor by radiative edge cooling
- special variants of the scenario (Quiescent H-mode, type II ELMs)
- ELM mitigation – pellet pacing or Resonant Magnetic Perturbations

Need to tackle NTM problem

- NTM suppression by Electron Cyclotron Current Drive demonstrated, but have to demonstrate that this can be used as reliable tool



- What is a 'tokamak (operational) scenario'?
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Advanced scenarios aim at stationary (transformerless) operation

- external CD has low efficiency (remember less than 0.1 A per W)
- internal bootstrap current high for high $j_{bs} \sim (r/R)^{1/2} \nabla p / B_{pol}$

$$\rightarrow f_{NI} \sim I_{bs}/I_p \sim p/B_{pol}^2 \sim \beta_{pol}$$

Recipe to obtain high bootstrap fraction:

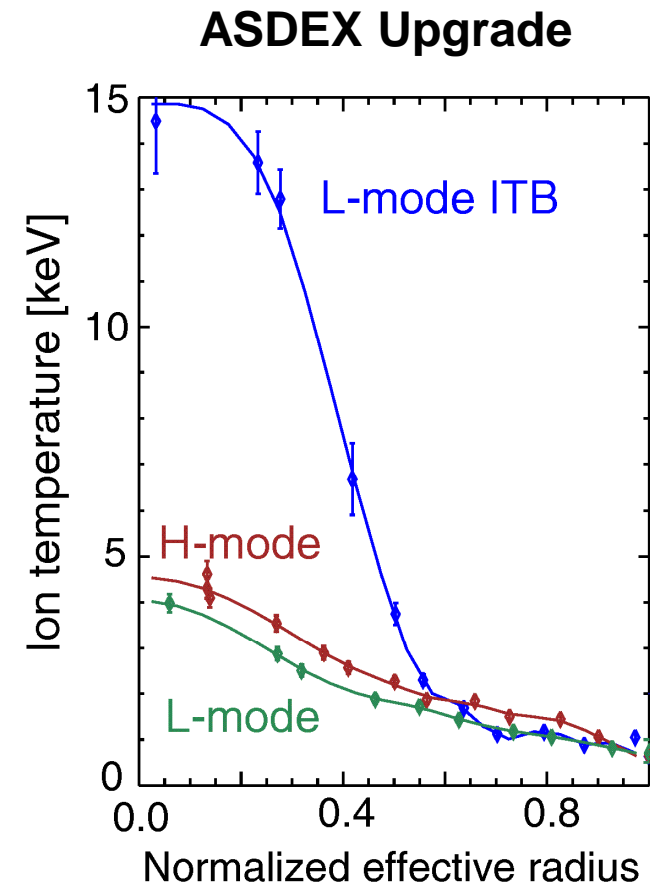
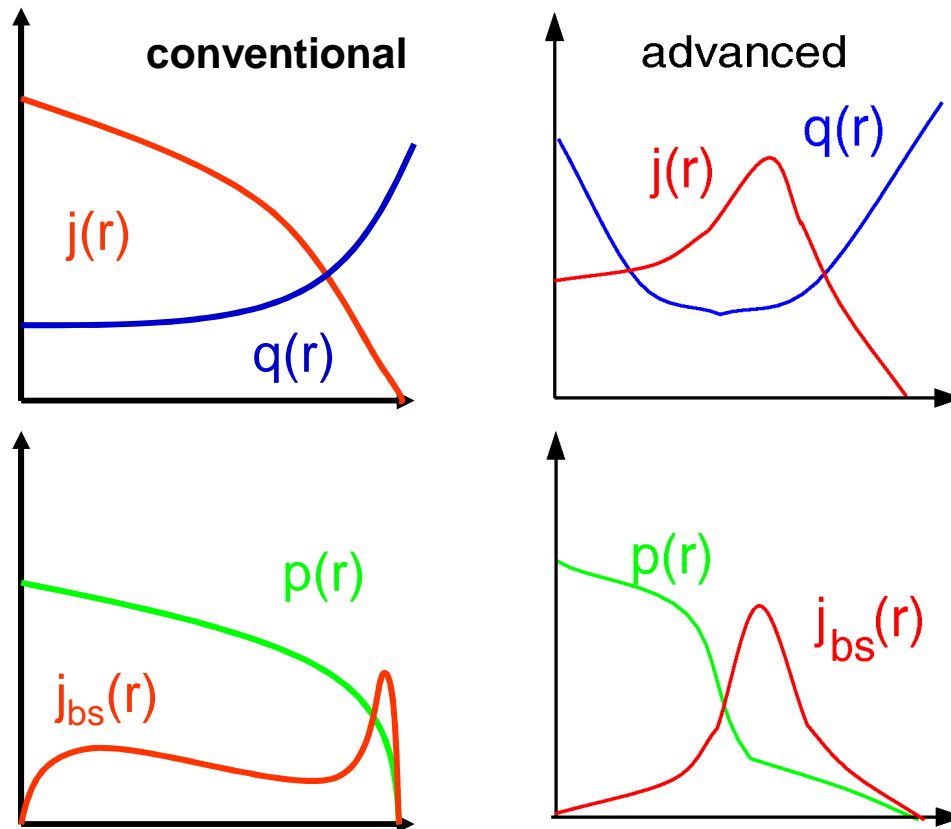
- low B_{pol} , i.e. high q – elevate or reverse q -profile ($q=(r/R)(B_{tor}/B_{pol})$)
- eliminates NTMs (reversed shear, no low resonant q -surfaces)
- high pressure where B_{pol} is low, i.e. peaked $p(r)$

Both recipes tend to make discharge ideal MHD (kink) unstable!

In addition, j_{bs} profile should be consistent with q -profile



Advanced tokamak – the problem of steady state



A self-consistent solution is theoretically possible

- reversing q -profile suppresses turbulence – internal transport barrier (ITB)
- large bootstrap current at mid-radius supports reversed q -profile



Broad current profile leads to low kink stability (low β -limit):

- can partly be cured by close conducting shell, but kink instability then grows on resistive time scale of wall (Resistive Wall Mode RWM)
- can be counteracted by helical coils, but this needs sophisticated feedback

Position of ITB and minimum of q -profile must be well aligned

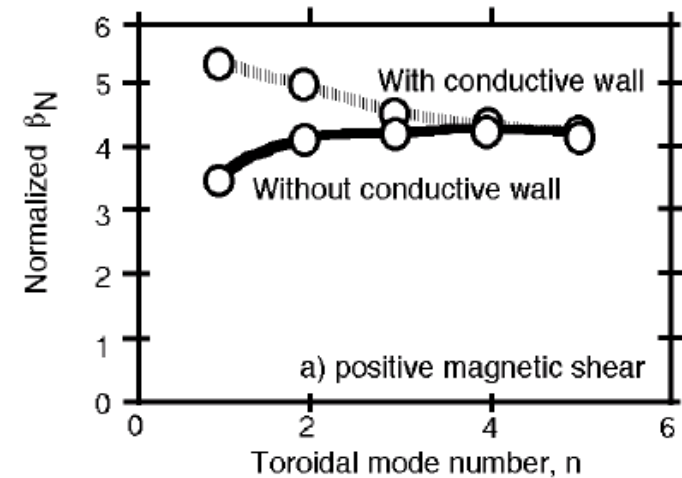
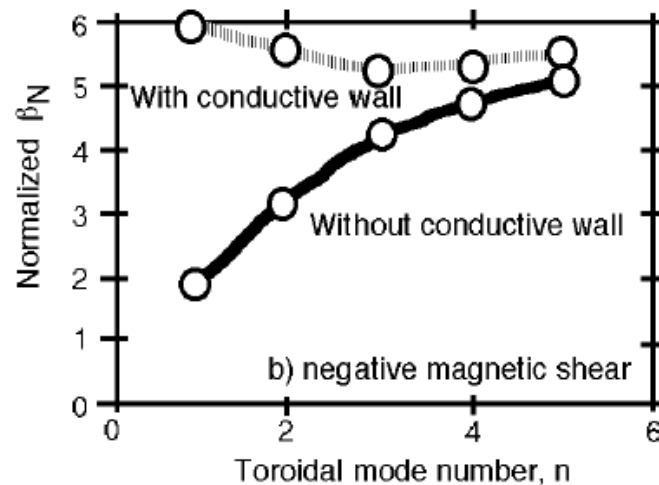
- needs active control of both $p(r)$ and $j(r)$ profiles – difficult with limited actuator set (and cross-coupling between the profiles)



Effect of a conducting shell on stability...



for conventional H-mode:
modest need and gain



for advanced ITB scenarios:
strong need and significant gain

Close-by conducting wall can stabilise external kink instability

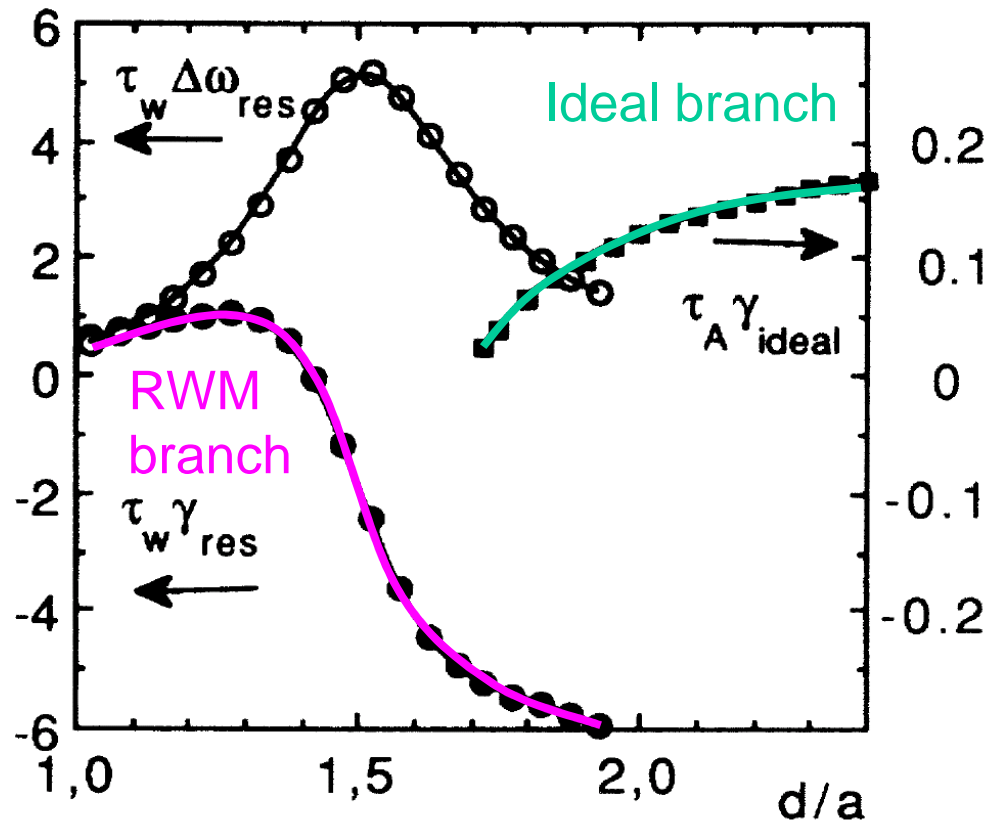
- usually not an issue for conventional scenario, but advanced scenarios prone to external kink due to broad current/peaked pressure profile



...but no wall is ideally conducting!



A. Bondeson et al., PRL 94



Ideal regime: Ideal kink stable if wall is close enough

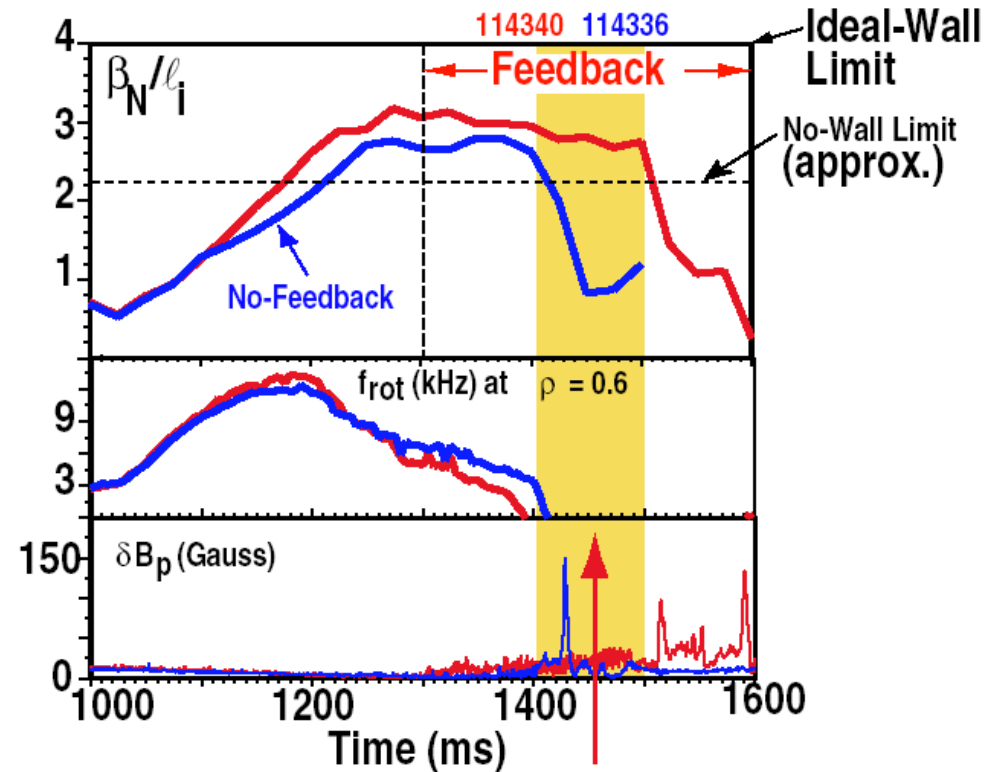
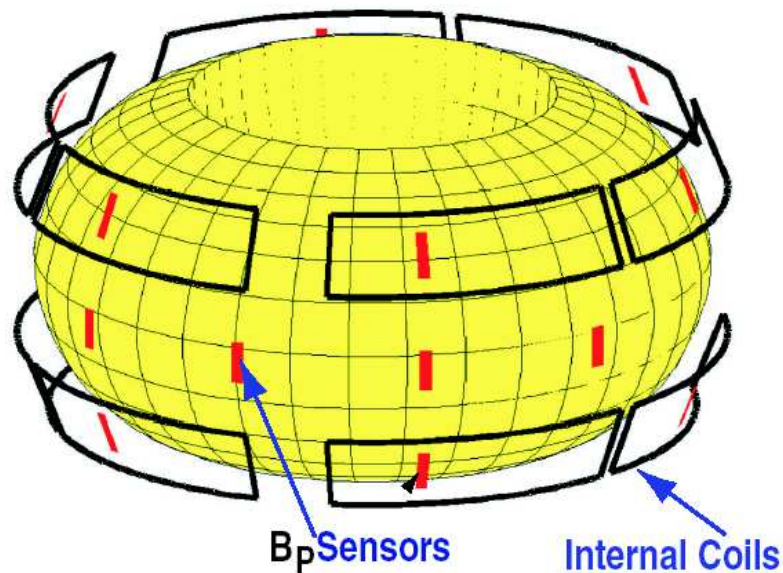
RWM regime: RWM is stable when slipping between mode and wall is large enough

When ideal kink is wall stabilised, RWM can grow on wall time scale

- rotation w.r.t. wall can stabilise the RWM if $\omega_{\text{rot}} \gg 1/\tau_w$
- balance between wall drag and (rotating) plasma drag on mode



RWM control by Resonant Magnetic Perturbations

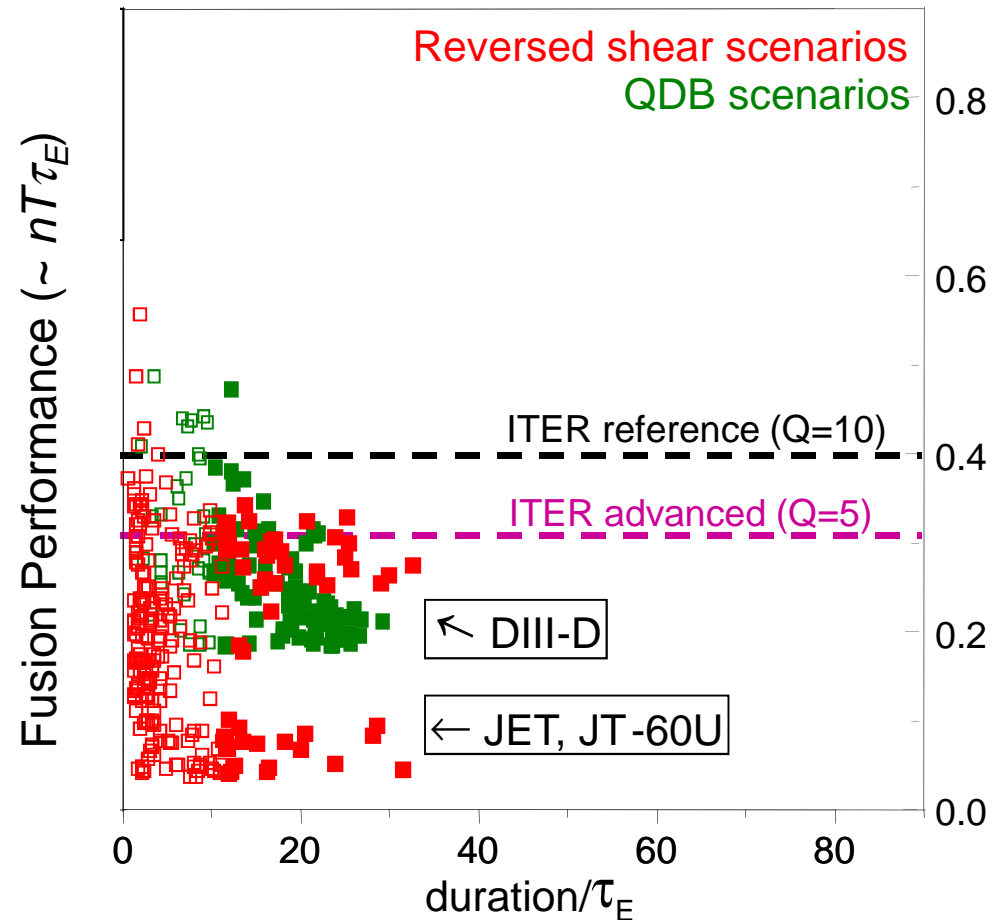


Feedback control using RMPs shows possibility to exceed no-wall β -limit

- rotation plays a strong role in this process and has to be understood better (ITER is predicted to have very low rotation)
- fast particle stabilisation can help substantially even at low rotation



Advanced Tokamak Stability is a tough Problem



Good performance can only be kept for several confinement times, not stationary on the current diffusion time (10 – 50 τ_E in these devices)



A compromise: the 'hybrid' scenario

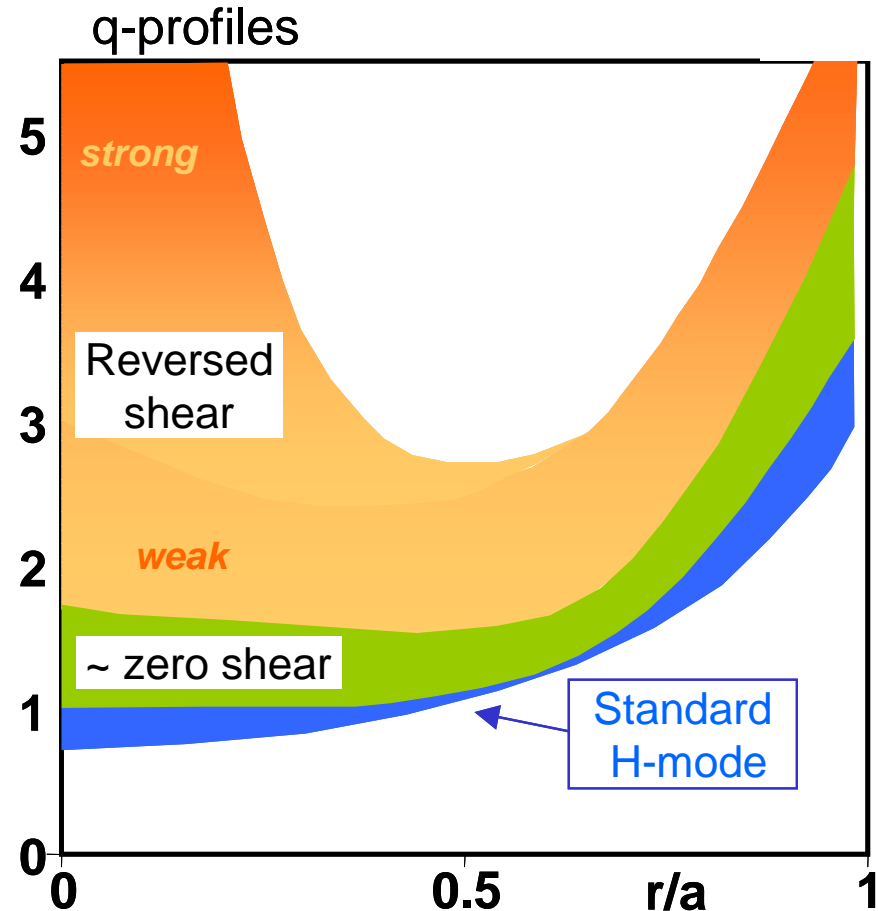


Reversed shear, ITB discharges

- + very large bootstrap fraction
- + steady state should be possible
- low β -limit (kink, infernal, RWM)
- delicate to operate

Zero shear, 'hybrid' discharges

- + higher β -limit (NTMs)
- + 'easy' to operate
- smaller bootstrap fraction
- have to elevate $q(0)$
(also avoids sawteeth, NTMs)



Hybrid operation aims at flat, elevated q-profile discharges with high $q(0)$

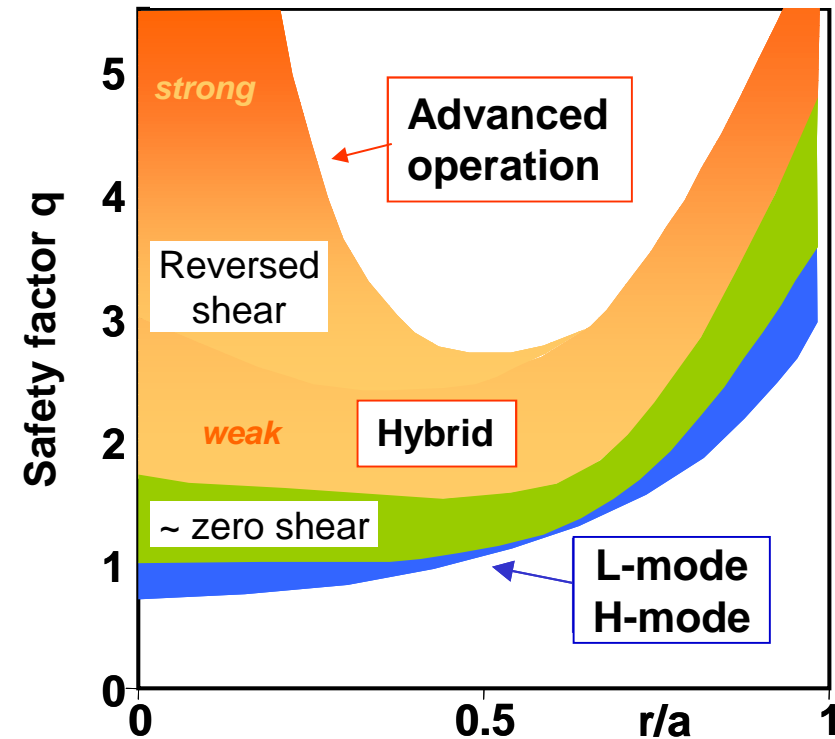
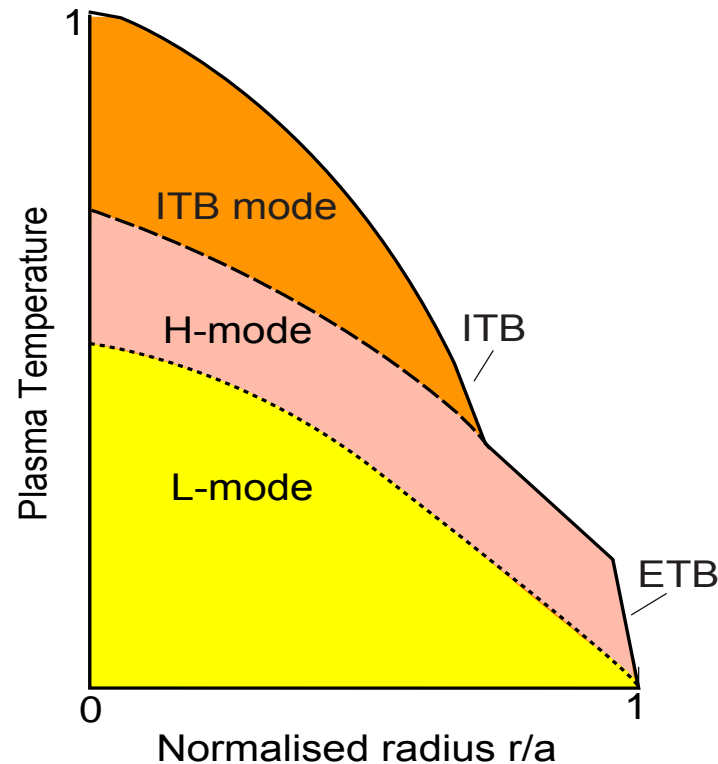
Not clear if this projects to steady state, but it will be very long pulse...



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Summary and Conclusions



A variety of tokamak operational scenarios exists

- L-mode: low performance, pulsed operation, no need for profile control
- H-mode: higher performance, pulsed operation, MHD control needed
- Advanced modes: higher performance, steady state, needs profile control



Summary and Conclusions



ITER aims at operation in conventional *and* advanced scenarios

- demonstrating $Q=10$ in conventional (conservative) operation scenarios
- demonstrating long pulse (steady state) operation in ,advanced‘ scenarios

Scenario:	Standard	Low q	Hybrid	Advanced
I_p [MA]	15	17	13.8	9
B_t [T]	5.3	5.3	5.3	5.18
β_N	1.8	2.2	1.9	3
P_{fus} [MW]	400	700	400	356
Q	10	20	5.4	6
t_{pulse} [s]	400	100	1000	3000

One mission of ITER and the accompanying programme is to develop and verify an operational scenario for DEMO

- DEMO scenario must be a point design (no longer an experiment)
- actuators even more limited (e.g. maximum of 2 H&CD methods)