

## Problem 7

### WARP-drive

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**Background** "WARP 5, Mr. Sulu".

What a dream. Travelling as fast as one wishes. But nature sets a boundary, the speed of light. Eventhough there might seem no chance for a 'WARP-drive', in the 90ies M. Alcubierre published a geometry of an Einsteinian space-time that would allow for transport of material objects over far spatial distances in arbitrary short amount of time. [1]

In this problem, we will study its metric and discuss the need for a certain arrangement of 'exotic matter' for such a space-time. We will work in two spatial dimensions. That eases the problem without loss of essential features of Alcubierre's ideas. In this problem, the signature convention on the metric is  $(-, +, +)$  and  $dx^0 = dt$ ,  $dx^1 = dx$ , and  $dx^2 = dy$ .

Let's begin by discussing the Alcubierre metric. It is constructed as follows

$$ds^2 = \sum_{\mu=0}^2 \sum_{\nu=0}^2 g_{\mu\nu} dx^\mu dx^\nu, \quad (7.1)$$

$$\{g_{00} = -(1 - \beta^2(t, x, y)), \quad g_{01} = g_{10} = \beta(t, x, y), \quad g_{11} = g_{22} = 1, \quad g_{\mu\nu} = 0 \text{ else}\}.$$

$\beta(t, x, y)$  models the 'WARP-bubble' in the  $(x, y)$ -plane for a given time.

$\beta(t, x, y) = 0$  aside the bubble.

*If you want to see an explicit expression of how  $\beta$  may look like, confer the very end of this problem, however, you don't need that information for this problem!*

#### PART I

This metric may appear weird at first sight. It has 'side elements' and  $g_{00}$  seems to change sign, depending on where one evaluates the metric in the parameterisation. So, let's check if it is a proper metric.

**a) [0.5 points]** Show that  $g$  is asymptotically flat, i. e. it converges to the Minkowski metric far aside the 'WARP-bubble' for any given time.

Hint: *The Minkowski metric reads  $\{\eta_{00} = -1, \quad \eta_{11} = \eta_{22} = 1, \quad \eta_{\mu\nu} = 0 \text{ else}\}$ .*

**b) [1 point]** Proof that its determinant reads

$$g = \det(g_{\mu\nu}) = -1, \quad (7.2)$$

and find the inverse metric  $g^{\mu\nu}$ . Note  $\sum_{\alpha=0}^2 g^{\mu\alpha} g_{\alpha\nu} = \delta_\nu^\mu$ , where  $\delta_\nu^\mu = 1$  for  $\mu = \nu$ , and 0 otherwise.

Hint: *you can use the fact, that  $g$  'looks like' a block diagonal matrix.*

Finally, the metric has to support time-like vector fields. A vector field  $\{u^\mu\}$  is time-like iff

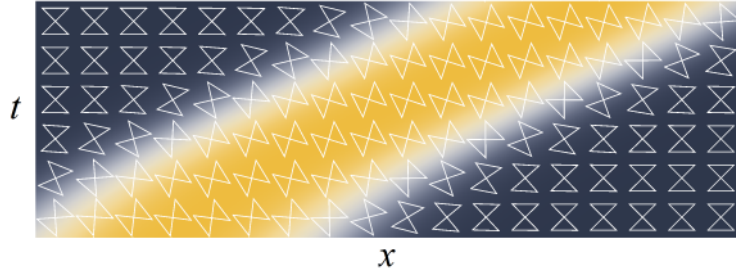
$$\sum_{\mu=0}^2 \sum_{\nu=0}^2 g_{\mu\nu} u^\mu u^\nu < 0. \quad (7.3)$$

c) [1 point] Check that

$$\{u^0 = 1, \quad u^1 = -\beta, \quad u^2 = 0\}, \quad (7.4)$$

is indeed such a time-like vector field. A trajectory to which  $u$  is tangential obeys the principle of special relativity to not exceed the speed of light locally.

In that context, describe **Figure 7.1**.



**Figure 7.1** Cross-section of Alcubierre space-time at  $y = 0$  for far away observer. The cones are the local light cones. [5]

## PART II

$\Gamma_{\nu\alpha\beta}$  are the Christoffel symbols (of first kind), which convey the gravitational field.

$$\Gamma_{\nu\alpha\beta} := \frac{1}{2} (-\partial_\nu g_{\alpha\beta} + \partial_\alpha g_{\beta\nu} + \partial_\beta g_{\nu\alpha}) \quad (7.5)$$

[2]. Einstein supposed that the energy momentum density tensor  $T^{\mu\nu}$  sources a non-trivial curvature to the geometry of our space-time. The component  $T^{00}$  in  $T^{\mu\nu}$  equals the energy density.

*Note that matter, as we know it, is represented by an energy density field that is strictly positive!*

Due to energy momentum conservation and symmetry constraints the simplest possible way to formulate such a relation between the energy momentum density tensor and curvature is

$$\kappa \cdot T^{\rho\sigma} = R^{\rho\sigma} - \frac{1}{2} g^{\rho\sigma} R, \quad (7.6)$$

where  $\kappa$  is a positive constant.  $R = \sum_{\tau=0}^2 \sum_{\zeta=0}^2 g_{\tau\zeta} R^{\tau\zeta}$  is the Ricci scalar, and  $R^{\tau\zeta}$  is the Ricci tensor. We

are going to use Einstein's sum convention, i. e. if an upper index and a lower index have the same letter and occur in a product of two tensors, it is automatically summed over:  $R = g_{\tau\zeta} R^{\tau\zeta}$ . The Ricci tensor can be constructed from the Riemann curvature tensor  $R_{\alpha\beta\mu\nu}$ , via

$$R^{\rho\sigma} = g^{\rho\beta} g^{\sigma\nu} R_{\beta\nu}, \quad R_{\beta\nu} = g^{\alpha\mu} R_{\alpha\beta\mu\nu}. \quad (7.7)$$

The Riemann curvature tensor contains all curvature information of our space-time and can be written as

$$R_{\alpha\beta\mu\nu} := \frac{1}{2} (\partial_\alpha \partial_\nu g_{\mu\beta} - \partial_\alpha \partial_\mu g_{\beta\nu} - \partial_\beta \partial_\nu g_{\mu\alpha} + \partial_\beta \partial_\mu g_{\alpha\nu}) + g^{\sigma\rho} (\Gamma_{\rho\alpha\nu} \Gamma_{\sigma\mu\beta} - \Gamma_{\rho\beta\nu} \Gamma_{\sigma\mu\alpha}). \quad (7.8)$$

Note exactly the order of indices! cf. [2].

Admittedly, this expression is overwhelming, thus this part of the problem is mostly about to find

an elegant way to not evaluating the most of these components and terms.

**d) [0.5 points]** Start out by showing

$$\Gamma_{\alpha\mu\nu} = \Gamma_{\alpha\nu\mu}, \quad (7.9)$$

$$R_{\alpha\beta\mu\nu} = R_{\mu\nu\alpha\beta} = -R_{\nu\mu\alpha\beta} = -R_{\mu\nu\beta\alpha}. \quad (7.10)$$

with using  $g_{\mu\nu} = g_{\nu\mu}$  and Eqn(7.5). Hint: *Symmetry relations of tensors transform to all coordinate frames. Only for this task you may go to a 'Riemann frame', in which the Christoffel symbols vanish!*

**e) [4 points]** Calculate all non-trivial Christoffel symbols.

The trivials are  $0 = \Gamma_{022} = \Gamma_{101} = \Gamma_{111} = \Gamma_{112} = \Gamma_{122} = \Gamma_{202} = \Gamma_{211} = \Gamma_{212} = \Gamma_{222}$ , and those that can be obtained via the symmetry relation.

*You do not have to validate these! As an intermediate result:  $\Gamma_{012} = \frac{1}{2}\partial_y\beta$ .*

**f) [0.5 points]** Only from these (anti-)symmetries you can work out what components of the Riemann tensor  $R_{\alpha\beta\mu\nu}$  are non-vanishing in general. Indicate them with a '×' in the following table :

$\begin{matrix} \mu\nu \\ \alpha\beta \end{matrix}$	00	01	02	11	12	22
00						
01						
02						
11						
12						
22						

For the remainder of the problem we want to work out what actually is the arrangement of matter we need, to make our space-time to become an Alcubierre space-time. To that end let's calculate  $T^{00}$ , and we will use some short cutting tricks, to not waste time on calculating Riemann components which we do not need.

**g) [1 point]** Write down  $R^{00}$  in terms of the explicit expression of the inverse metric and keep the components of  $R_{\beta\nu}$  as they come, like:

$$R^{00} = g^{0\beta}g^{0\nu}R_{\beta\nu} = g^{00}g^{00}R_{00} + \dots = (-1)(-1)R_{00} + \dots \quad (7.11)$$

Proceed in the same way for the Ricci scalar  $R$ .

**h) [1.5 points]** Now, calculate the energy density  $\kappa T^{00} = R^{00} - \frac{1}{2}g^{00}R$  by expanding all occurring Ricci-tensor components in terms of Riemann-tensor components, and show that this energy density is non-positive! **Figure 7.2** is an image of the resulting energy density.

Hint: *Do not calculate the Riemann components until the very end. You should find  $\kappa T^{00} = R_{2121}$ .*

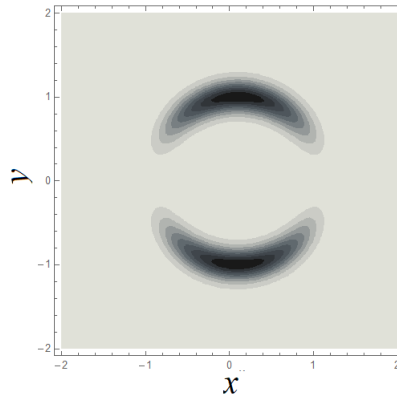
**A comment on how  $\beta(t, x, y)$  might look like.**

$$\beta(t, x, y) = -v_0 \cdot f(t, x, y), \quad (7.12)$$

where  $v_0 = \text{const.}$   $f$  is a smoothed Heaviside function on a circular domain (of radius  $r_0$ ) in the  $(x, y)$ -plane, and has a value of 1 inside the domain and 0 outside, with a smooth but narrow transition zone (of width  $w$ ) between these two values, so,  $f$  is a continuously differentiable function to at least second order. Its center is shifted along the  $x$ -axis over time.

$$f(t, x, y) = \Theta_w \left( \frac{r_0 - \sqrt{(x - v_0 \cdot t)^2 + y^2}}{w} \right). \quad (7.13)$$

$x, y$  and  $t$  is the choice of a global parameterisation of Alcubierre's space-time, used by an observer that is spatially far aside the 'WARP-bubble' for any given time.



**Figure 7.2** Arrangement of 'exotic matter' in  $(x, y)$ -plane at  $t = 0$  for  $w = 0.25$ . The darker the gray tone, the higher the density of 'exotic matter'. In 3D, a space ship being placed in the WARP-bubble given by  $f$ , would be surrounded by a 'donut' of 'exotic matter' with the symmetry axis being *along the x-axis!* [5]